

# Capacitive effects on quantitative dopant profiling with scanned electrostatic force microscopes

Todd Hochwitz, Albert K. Henning, Chris Levey, and Charles Daglian  
Dartmouth College, Hanover, New Hampshire 03755

James Slinkman  
IBM Microelectronics, Essex Junction, Vermont 05452

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A force-based scanning Kelvin probe microscope has been applied to the problem of dopant profiling in silicon. Initial data analysis assumed the detected electrostatic force couples the sample and only the tip at the end of a force sensing cantilever. Attempts to compare measurements quantitatively against device structures with this simple model failed. A significant contribution arises from the electrostatic force between the sample and the entire cantilever, which depends strongly upon the relative size of the tip, cantilever, and lateral inhomogeneities in the surface topography and material composition of the sample. Actual and simulated measurements which demonstrate the characteristic signature of this effect are presented. © 1996 American Vacuum Society.

## I. INTRODUCTION

Qualitative measurements taken with a combination of an atomic force microscope (AFM) and scanning Kelvin probe force microscope (SKPM) on simple implanted structures in silicon have been reported by several researchers.<sup>1–7</sup> It is reported in Ref. 6 that initial attempts to quantify the measurements against theoretical models have not been completely successful. Further work on the samples described in Ref. 6 is reported here.

One of the most significant reasons for the reported discrepancy between measured surface potential profiles obtained via the SKPM and profiles predicted by theoretical models is the interaction between the cantilever of the force sensing probe, in addition to the tip, and the sample surface. The standard electrostatic force model is presented and extended to show this discrepancy. Simulation and experimental results which illustrate this interaction are also presented.

## II. MODELS

A distinction is made between a cantilever, a tip, and a probe. For this work the cantilever is the section of the force sensor which undergoes mechanical deflection. The tip is assumed to be physically attached to the cantilever. It is the location of primary force interaction with a sample under investigation. The probe consists of both the cantilever and the tip.

The mechanical deflection of a cantilever due to an electrostatic force between the probe and sample has been examined by several researchers and treated as a capacitive system.<sup>1,8–11</sup> The force on the probe due to a potential difference between the probe and sample may be represented as<sup>2</sup>

$$F_e = -\frac{1}{2}V^2C', \quad (1)$$

$$C' = \left(\frac{C_{\text{eff}}}{C_{\text{air}}}\right)^2 \frac{\partial C_{\text{air}}}{\partial z}, \quad (2)$$

where  $V$  is the total potential difference,  $z$  is the spatial distance,  $C_{\text{air}}$  is the capacitance, and  $\partial C_{\text{air}}/\partial z$  is the spatial derivative of the capacitance between the probe and sample;  $C_{\text{eff}}$  is the total capacitive load in series with the system, including contributions from the probe/sample gap, dielectric materials on probe or sample surfaces, and bulk  $p/n$  junctions.

The above equations are most applicable to a system where the sample and probe are metals. In this case the only spacing-dependent capacitance between the probe and the sample is across the air gap. If either the probe or the sample is a nondegenerately doped semiconductor, however, the total charge contained within a volume near the surface may also vary as the probe/sample spacing changes, resulting in a different form of Eq. (2). In addition, filled surface states may lead to a different electrostatic force than predicted by Eq. (1).

Theoretical analysis indicates the cantilever deflection is directly proportional to the magnitude of this force.<sup>12</sup> Thus, variations in the electrical nature of the probe/sample system can be examined by detecting and properly adjusting for changes in the mechanical deflection of the cantilever.

With the scanning electrostatic force microscope (SEFM) the potential difference  $V$  between the probe and sample can be comprised of a dc signal and/or one or more ac signals. With the SKPM model presented in this article, the potential difference  $V$  consists of one dc component and one ac component, which may be represented as follows:

$$V = V(t) = V_{\text{dc}} - \frac{\Delta\Phi}{q} + V_{\text{ac}} \sin(\omega t), \quad (3)$$

where  $V_{\text{dc}}$  is the magnitude of the dc component applied by any external electronics,  $\Delta\Phi$  is the work function difference (in eV) of the materials [and  $\Delta\Phi/q$  is the contact potential difference (CPD)],  $V_{\text{ac}}$  is the peak magnitude of the ac com-

ponent of the externally applied potential between the probe and the sample, and  $\omega$  is the frequency at which the ac voltage is oscillating.

If  $V$  in Eq. (1) is replaced by the form given in Eq. (3) and expanded, one obtains

$$F_e = \frac{1}{2} \left( \left( V_{dc} - \frac{\Delta\Phi}{q} \right)^2 + \frac{1}{2} V_{ac}^2 \right) C' \quad (4)$$

$$+ C' \left( V_{dc} - \frac{\Delta\Phi}{q} \right) V_{ac} \sin(\omega t) \quad (5)$$

$$+ \frac{1}{4} C' V_{ac}^2 \cos(2\omega t). \quad (6)$$

The negative sign from Eq. (1) has been removed since the force is always attractive under the circumstances investigated during this study. An interferometer is used to monitor cantilever deflections at the frequencies of  $\omega$  and  $2\omega$ . The deflection at  $2\omega$ , which is proportional to Eq. (6), yields sample capacitance information. The signal detected at  $\omega$ , which is proportional to Eq. (5), is used to minimize the electrostatic force. Since the capacitive and ac terms of Eq. (5) are nonzero, the primary method to minimize cantilever deflection at  $\omega$  is to adjust  $V_{dc}$  until it equals the CPD between sample and probe.

Ideally, this feedback signal is used to estimate spatial variations in the sample's surface work function (WF). In practice the feedback signal is equal to the CPD only for samples with surface features significantly larger than the entire probe. Our experience shows it is necessary to consider the interaction of all the electrically conductive components of the force sensing probe and sample surface, as suggested by other researchers.<sup>9,13,14</sup>

Since the sample surface varies under the probe, the value of  $V_{dc}$  which minimizes the cantilever deflection becomes a weighted average of the CPD and capacitive components between the probe and the sample surface. This situation is similar to the stray capacitance problems associated with the displacement current-based Kelvin probes.<sup>15–19</sup> The impact of the errors is very different in the two techniques, so a direct application of prior analysis is not appropriate.

An example of this stray capacitance for the force-based SKPM is shown in Fig. 1. It shows the profile of a micro-fabricated silicon force sensing probe<sup>20</sup> used in our system as it scans at a height of  $Z_0$  over a rectangular surface feature with vertical dimension  $\Delta S$ . In general, the total capacitance between the probe and sample would be an integral computed over the entire electrically active surface.<sup>14</sup> For this work the capacitive components will be treated discretely to simplify the discussion and numerical analysis.

The total load shown in Fig. 1 between the electrical connections to the sample and probe is modeled as three parallel capacitances. One component is between the end of the cantilever and the sample, the second component is between the tip and the feature of interest, and the third is between the bulk of the cantilever and the sample. Since these components are in parallel, the total effective capacitance is a sum of the individual capacitors. These capacitances are also in

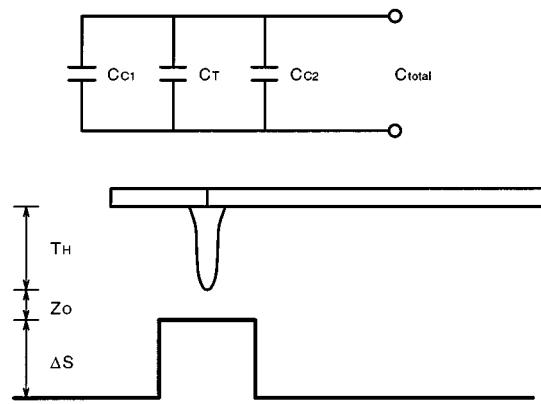


FIG. 1. Schematic representation of the probe/sample subsystem and the equivalent capacitance circuit. The desirable capacitance is between the tip and the sample surface ( $C_T$ ). The remaining cantilever and sample overlap yields stray capacitance ( $C_{C1}$  and  $C_{C2}$ ). The tip is kept within tens of nanometers from the sample surface ( $Z_0$ ) while the cantilever is several microns above the tip ( $T_H$ ).

series to any additional capacitance between the electrical connections and the air gap, which for now is assumed to be constant.

The entire system can be broken down further and modeled as several very small capacitors in parallel, with the effective capacitance given as

$$C_{total} = \sum_i C_i, \quad (7)$$

where  $C_i$  is the contribution from each capacitive term. The derivative of  $C_{total}$  taken with respect to the vertical displacement between the sample and probe becomes

$$\frac{\partial C_{total}}{\partial z} = \sum_i \frac{\partial C_i}{\partial z}. \quad (8)$$

The nature of Eq. (5) is thus changed. The electrical signal  $S_\omega$  obtained from a lock-in amplifier becomes

$$S_\omega \propto V_{ac} \sum_i \frac{\partial C_i}{\partial z} \left( V_{dc} - \frac{\Delta\Phi_i}{q} \right). \quad (9)$$

The applied voltage  $V_{dc}$  needed to minimize the deflection of the cantilever at  $\omega$  is derived as follows:

$$\sum_i \frac{\partial C_i}{\partial z} \left( V_{dc} - \frac{\Delta\Phi_i}{q} \right) = 0, \quad (10)$$

$$V_{dc} = \frac{\sum_i (\partial C_i / \partial z) (\Delta\Phi_i / q)}{\sum_i (\partial C_i / \partial z)}. \quad (11)$$

In addition, Eq. (6) will also be modified. The capacitive term will be replaced with the summation from the denominator of Eq. (11). Therefore, the capacitance signal will play an important role in extracting quantitative CPD results from the  $V_{dc}$  required to minimize the electrostatic force.

This description is a simplified version of the contribution from the various capacitive components to the detected can-

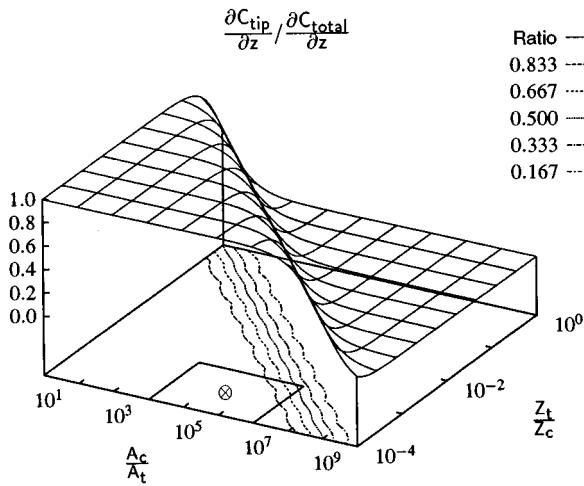


FIG. 2. A surface plot showing the relative contribution of the tip/sample capacitance compared to the total probe/sample capacitance as functions of the area and sample spacing over a topographically flat surface ( $\Delta S=0$ ).  $A_c/A_t$  is the ratio of the cantilever area to the tip area and  $Z_t/Z_c$  is the ratio of the tip/sample spacing to the cantilever/sample spacing.  $Z_t$  corresponds to  $Z_0$ , and  $Z_c$  to  $Z_0+T_H$ , from Fig. 1.

tilever deflection. To model the cantilever deflection more correctly, one would need to perform a self-consistent calculation of the equations of motion with the contributions from the entire probe. As described earlier, the summation is an approximation to the integral required to obtain the exact relationship between the probe voltage and material parameters.

The simple analysis indicates the feedback voltage depends upon the lateral extents of the surface topography, dielectric properties, and CPD over the electrically active area between the probe and the sample. In order to obtain high resolution, it is necessary for most of the contribution to come from the interaction of the tip/sample and not the cantilever/sample. Figure 2 shows a plot of the following relationship:

$$\frac{\partial C_{\text{tip}}/\partial z}{\partial C_{\text{tip}}/\partial z + \partial C_{\text{cantilever}}/\partial z} \quad (12)$$

as a function of physical cantilever/tip dimensions and probe/sample spacing assuming a topographically flat perfect conductor, where  $\partial C_{\text{tip}}/\partial z$  is the effective tip/sample capacitance and  $\partial C_{\text{cantilever}}/\partial z$  is the effective cantilever/sample capacitance. The capacitance was treated as a parallel plate, so the curve shown is applicable only to systems that may be described as being proportional to an effective surface area and inversely proportional to the electrode spacing. The shape of the curve for other systems will vary, but the trends should remain similar.

The silicon probes used are fabricated to be  $\approx 100 \mu\text{m}$  long  $\times 20 \mu\text{m}$  wide. The tip has a radius of curvature of  $\approx 20 \text{ nm}$  when sharp, and  $\approx 0.3\text{--}0.5 \mu\text{m}$  when dulled after many scans or due to tip crashes. The normalized area of cantilever to tip lies in the range of  $10^3\text{--}10^6$ . The force-gradient control

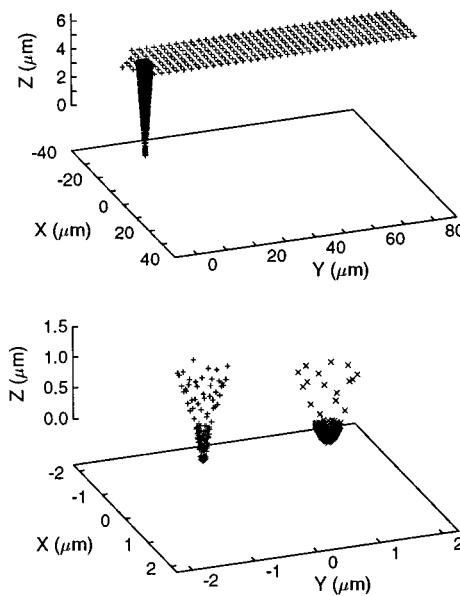


FIG. 3. Plots showing the elements used to model the cantilever, a sharp (+) tip, and a dull (x) tip of a force sensing probe. The cantilever is  $20 \mu\text{m}$  wide  $\times 100 \mu\text{m}$  long. The sharp tip has a radius of curvature of  $\approx 20 \text{ nm}$ , and the dull tip has a radius of curvature of  $\approx 0.25 \mu\text{m}$ .

loop of our system is typically set to keep the tip  $\approx 5 \text{ nm}$  from the surface ( $Z_0$  in Fig. 1), and the tip length is specified to be in the range of  $6.5\text{--}7.0 \mu\text{m}$  ( $T_H$  in Fig. 1). The tip/sample to cantilever/sample spacing ratio is less than  $10^{-3}$ , which puts our operating area within the marked section of the plot ( $\otimes$ ).

More than 99% of the capacitance spatial derivative comes from the tip/sample for dull tips. For sharp tips this drops to only 63%. If one attempts to increase resolution by fabricating finer tips, or if there is significant relief in the surface topography changing the spacing ratio, the tip capacitance derivative ratio drops below 50%. Even if the feature under investigation is only as large as the effective area of the tip, the feedback voltage will not be the same as the material CPD unless the cantilever is also above material with a similar WF.

Our system is qualitatively consistent with this behavior. Dull tips yield much more contrast at a loss in lateral resolution. Measurements taken with tungsten probes yield more contrast due to the larger effective tip area and a value of  $T_H$  in excess of  $150 \mu\text{m}$ .

### III. SIMULATION RESULTS

To determine the impact of this modified formulation upon the expected measurements, we computed Eq. (11) for various types of probes and samples. The probe was divided into several elements, as shown in Fig. 3. For each of these elements, the corresponding  $C_i$ ,  $\partial C_i/\partial z$ , and  $\Delta\Phi_i/q$  were computed between the  $i$ th element of the probe and the nearest point on the model of the sample's surface. The capaci-

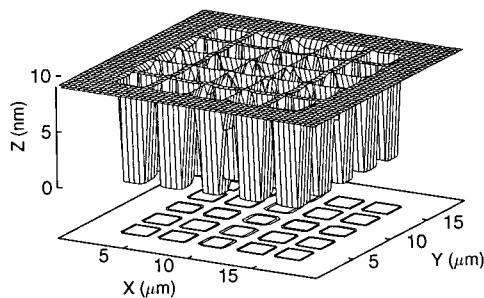


FIG. 4. Plot of the idealized surface topography for the model of the dopant contacts described in Ref. 6. The contact holes are arranged in  $5 \times 5$  arrays, with a width of  $2 \mu\text{m}$  for each contact hole. The depth of each contact is 9 nm.

tance between the  $i$ th element and the surface was treated as a parallel plate in order to speed computation.

A finite element model (FEM) was used to represent the surface characteristics of the sample. The values of the surface height, material WF, and effective substrate capacitance are stored for each point in the FEM mesh. The resulting probe voltage of Eq. (11) is computed for each point in the mesh. For every probe element that does not lie on a mesh node, or when a computed value is desired for the surface at a point that does not lie on a mesh node, the values are computed via linear interpolation.

Figure 4 shows a perspective plot of the topography associated with the doped structures described in Ref. 6. There is an anomalous structure in the figure introduced by the plotting package. It may be seen in the surface plot in the region  $10 \mu\text{m} < X < 15 \mu\text{m}$  and  $0 \mu\text{m} < Y < 5 \mu\text{m}$ , but not in the contour plot. The structure is *not* present in the mesh data used for the simulations. Figure 5 shows profiles of the effective substrate capacitance and WF for a slice taken through the middle of the structures. The values were representative of SIMS profiles and theoretical models for this particular sample.

A simulation of the scan was then performed with four types of probes. Two of the probes had sharp tips, and two had dull tips. For each pair, one probe was all silicon and the other had a silicon tip and gold cantilever. No simulation of an all gold probe was performed because our experience indicates that the gold tends to flake off, or the tip dulls, after extended use.

Figure 6 shows the simulated surface potential profiles for each of these four probes. The profiles are taken perpendicular to the direction of the simulated scan. The bulk of the cantilever was directed to the left. During the simulated scan the tip of the probe reaches the structures before the body of the cantilever, which mirrors the experimental procedure.

As predicted by our model, and consistent with the behavior of our SKPM, more contrast is indeed predicted with a dull tip compared to a sharp tip. There is slope in the predicted signal, which is related to how much of the cantilever resides over the structures. The contrast is larger for the right

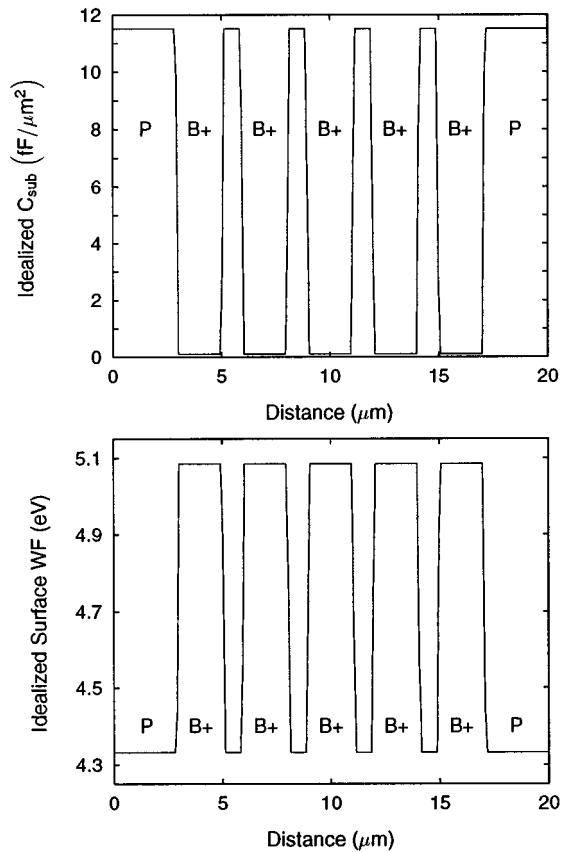


FIG. 5. Profiles of the idealized effective surface capacitance and work function used in the model. The substrate is assumed to have 3 nm of oxide present. The capacitance within the contact holes is a much lower value due to the junction capacitance of the implanted  $p/n$  depletion region being in series with the surface oxide. The substrate was  $n$ -type silicon, doped with phosphorus to a density of  $10^{14} \text{ cm}^{-3}$ , with a work function of 4.33 eV. The contact holes were doped with boron to a density of  $10^{19} \text{ cm}^{-3}$ , with a work function of 5.09 eV.

structure because the cantilever is interacting with all five implant regions. There are also effects seen in the profile due to the physical edges of the surface.

One interesting item is the change in contrast for a sharp Si/Au probe, where the predicted surface potential for the implanted structures is actually less than that predicted by theory (compare to Fig. 5). In several experimental situations a contrast change was found in the SKPM measurements. This change depends upon the settings of the force-gradient control loop when using coated probes. The contrast inversion may be attributed to changes in the relative contributions of the Si/sample and Au/sample CPDs to the deflection of the cantilever.

Future software revisions will include self-consistent solutions to the equations of motion and electrostatic forces. This change will allow circumvention of the approximations present in Eq. (11).

#### IV. EXPERIMENTAL RESULTS

Presented here are measurements made with the SKPM on the samples described in Ref. 6. The measurements were

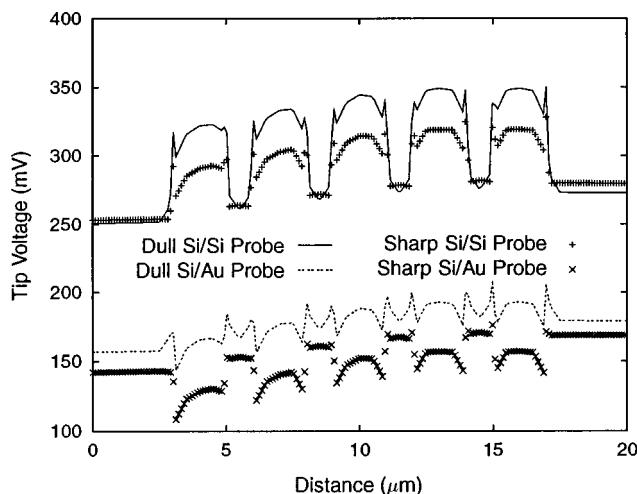


FIG. 6. Profiles of the simulated probe voltage. The profiles are taken through the five contact holes in the middle of the array, and perpendicular to the direction of the simulated scan. The leftmost points correspond to a scan point lying on the first scan line, and the rightmost points correspond to a scan point lying on the final scan line. The results depend strongly upon how much of the tip and cantilever lie above the array, in addition to the material properties of the tip and cantilever.

performed upon structures in a different area of the test die, and so are not identical to previously published results. The test die contains several test structures grouped together in different sizes.

The measurements reported in Ref. 6 were taken near a center group, using a probe with a tip radius of curvature less than 50 nm. The measurements reported here are taken near an edge group, using a probe with a tip radius of curvature around 300 nm. As shown in Sec. III, when the cantilever covers a larger number of the implanted structures the signal better approximates the true CPD. The previously published measurements show a wider signal swing than the ones in this article, consistent with the model presented earlier.

Figure 7 shows the actual measurement profiles obtained with the SKPM. To remain consistent with the simulation results, the lines are taken perpendicular to the direction of the scan. The leftmost signal peaks correspond to implants that are in the interior of an array, and the rightmost peaks are on the edge of the implanted array.

The absolute value of the probe voltage does not match the simulated value. This difference is likely due to the condition of either the sample or probe surface. The signal swing of the simulation was  $\approx 75$  mV, whereas for the measured data it was  $\approx 110$  mV. Possible reasons for this discrepancy are uncertainty in the dopant concentrations near the surface of either the sample or probe, incorrect  $p/n$  junction capacitance or oxide thickness, incorrect probe/sample spacing, incorrect tip size, coarse discretization of the sample mesh used in the simulations, an invalid assumption that only the capacitance across the air gap depends upon the tip/sample spacing, ignoring the presence of surface states or a thin water layer, ignoring the interaction of the laser with the silicon probe, off-axis orientation of the probe and sample

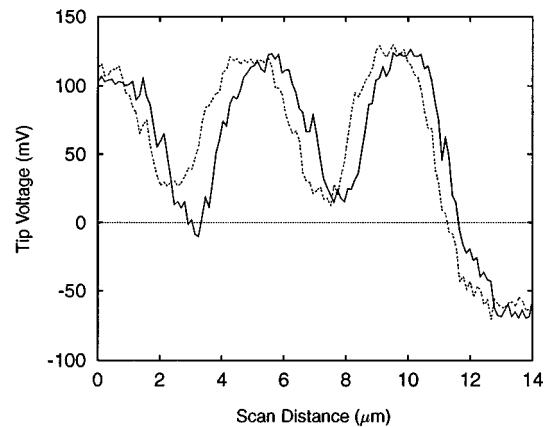


FIG. 7. Profiles of the actual probe voltage from measurements performed on a fabricated structure of contact holes. The profiles are taken perpendicular to the scan direction in order to match the conditions of the profile shown in Fig. 6. The peaks exhibit a slope similar to that shown in Fig. 6, and the edges show a similar discontinuity. The discontinuity in the simulation is stronger due to the approximations used in the model.

surface, and/or interaction with structures on the sample surface that were not included in the FEM mesh.

Though the results are not quantitatively exact, they are qualitatively similar. The measured data show a similar slope in the signal as more of the cantilever covers the implant areas. Discontinuities are seen along the edges of the implant areas (they are not as strong as in the simulation results due to the method of numerical analysis), and the crowns of the features are rounded.

## V. CONCLUSION

A simple extension of the electrostatic force model leads to theoretical results closer to the actual surface potential profiles measured with a force-based SKPM than is obtained with a traditional model. The discussion here has concentrated upon the analysis of previously published measurements.<sup>6</sup> The analysis has also been applied to other simple doped structures, including step- and stripelike features.

The role of the cantilever cannot be ignored in the quantitative analysis of data obtained with an SEFM, especially where claims about the lateral resolution are made. The greatest impact will occur in dc measurements (the Kelvin mode), since the WF is essentially a dc signal with an infinite lateral extent. There will be less impact on the ac measurements<sup>21–23</sup> if the separation of signal lines carrying the ac signals is much larger than the cantilever size. When the cantilever lies over two different lines carrying signals at the same ac frequency, or when it lies over a significant portion of a single line, the probe signal will not reflect the actual probe/surface signal (amplitude and phase). The presence of the extra capacitance will also have an impact upon the amount of current drawn through the system, which may influence the sample under study.

More quantitative statements about the measurements will require further deconvolution of the data. Our future research includes the formulation of a robust technique to extract the actual CPD and capacitance variations from the electrostatic and capacitance measurements obtained with the SKPM.

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